

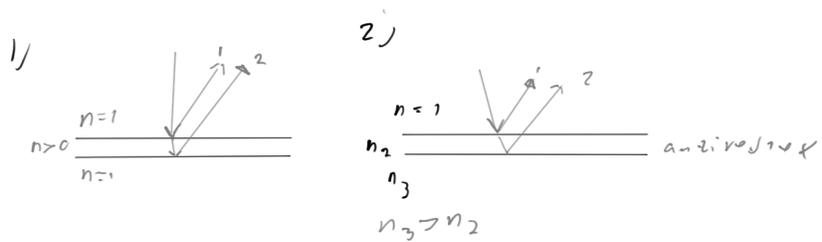
# Lecture notes

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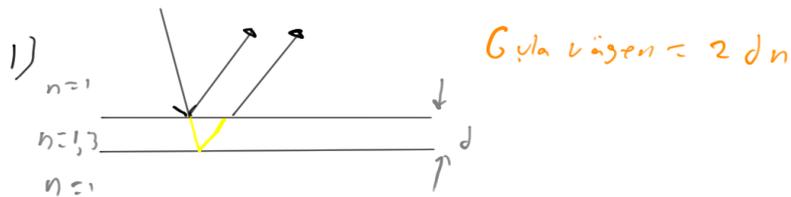
5 December 2016

1. Interferens i tunna filmer
2. Diffraktion
3. Upplösningförmåga
4. Rotations rörelser

## Reflektion i tunna filmer



### Fall 1



Fas skillnad:  $2\pi \frac{2dn}{\lambda} \pm \pi$

max:  $2\pi \frac{2dn}{\lambda} \pm \pi = m2\pi$

$$2d = \left(m \pm \frac{1}{2}\right) \lambda$$

min:  $2dn = m\lambda$

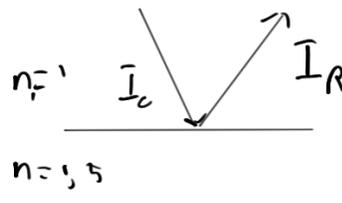
Fall 2

$$\text{max: } 2dn = m\lambda \quad m = 1, 2, 3, \dots$$

$$\text{min: } 2dn = \left(m + \frac{1}{2}\right)\lambda \quad m = 0, 1, 2, 3, \dots$$

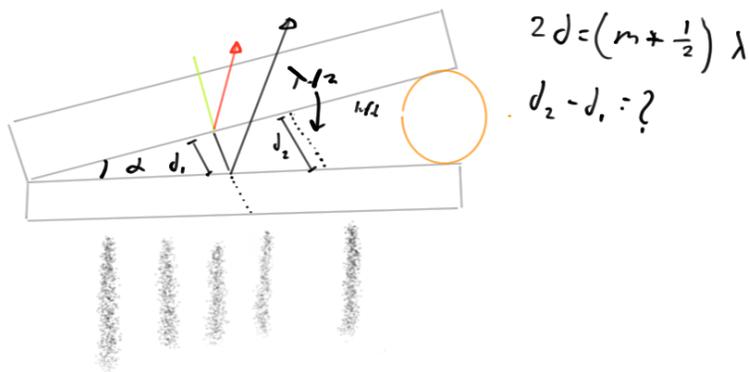
minsta tjocklek som ger min:

$$d = \frac{\lambda}{4n}$$



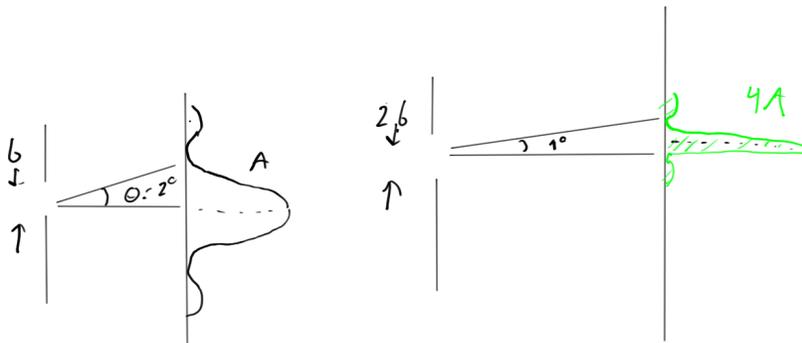
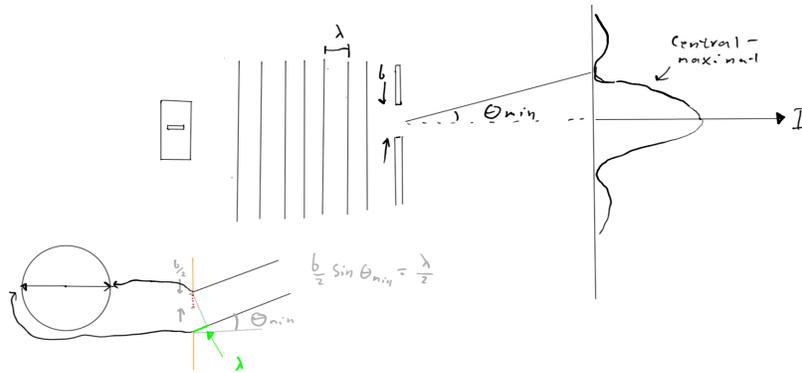
$$I_R = I_i \left( \frac{n_2 - n_1}{n_2 + n_1} \right)^2$$

Antireflexbehandling av glasögon:

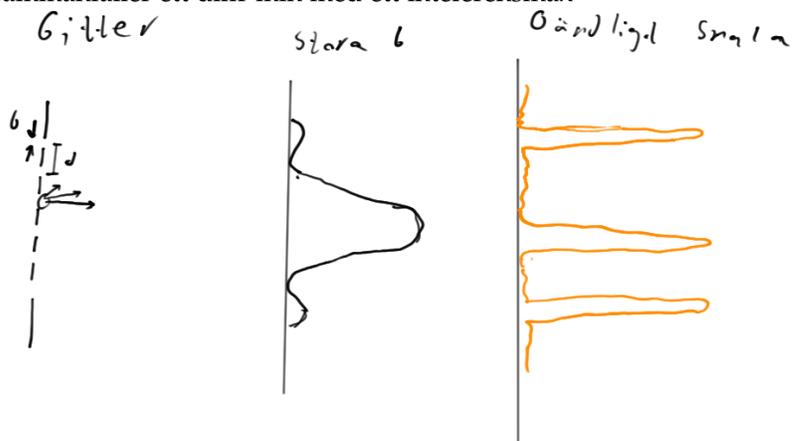


Diffraction

$$\text{min: } b \sin \theta_{\min} = \lambda \quad | \quad \text{Young, max: } d \sin \theta = \lambda$$



När sammanfaller ett diffr min med ett interferensmax

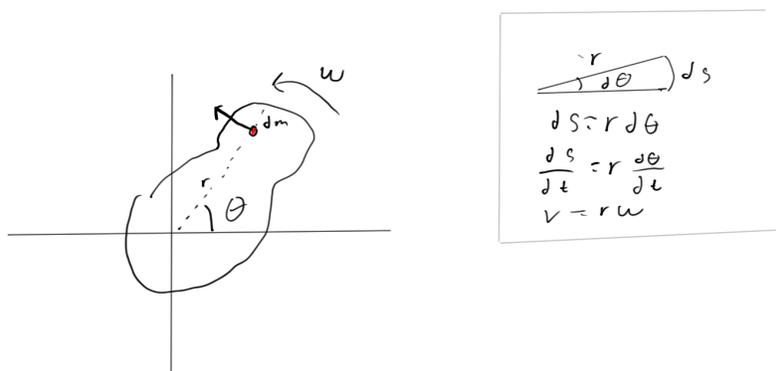
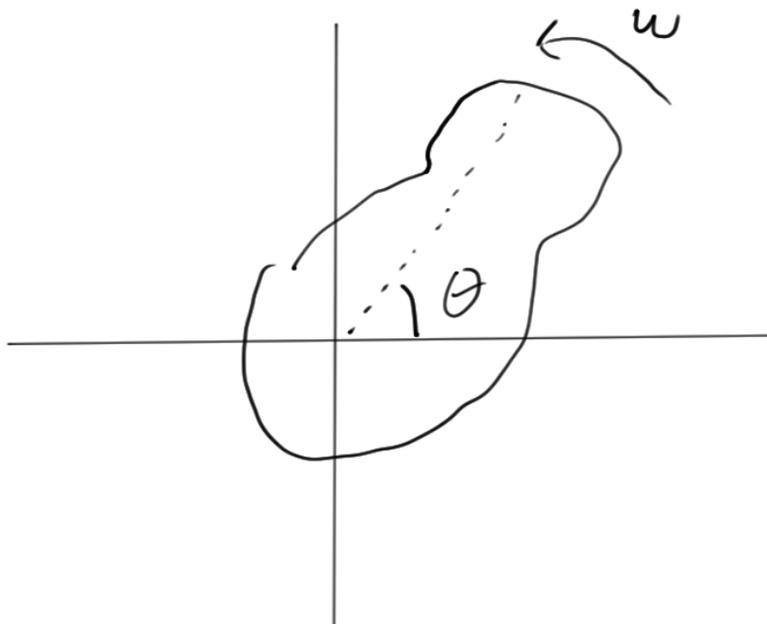


$$b \sin \theta = \lambda, d \sin \theta = m\lambda \rightarrow \frac{d}{b} = m$$



## Rotationsrörelse

## Rotationskinematik



läge:  $\theta$  jfr.  $x$

$$\text{hastighet: } \omega = \frac{d\theta}{dt} \vec{v}$$

$$\text{acceleration: } \alpha = \frac{d\omega}{dt} = \frac{d^2\theta}{dt^2} \vec{a}$$

$$\theta_f - \theta_i = \omega_0 t + \frac{1}{2} \alpha t^2$$

$$\omega_f^2 - \omega_i^2 = 2\alpha \Delta\theta \text{ Rotations energi:}$$

$$dK_R = \frac{1}{2} dm \cdot v^2$$

$$v = \omega r$$

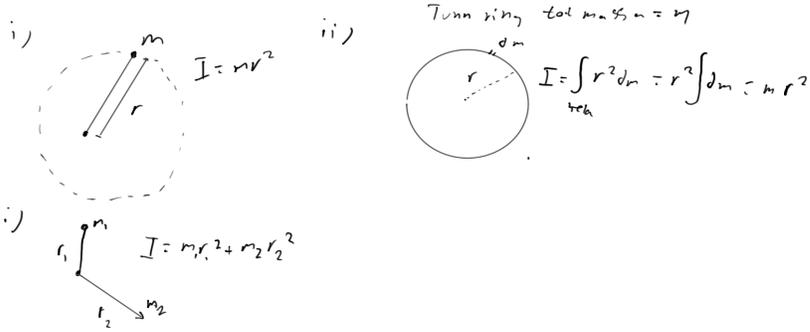
Hela rotations energin:

$$K_R = \int_{\text{hela}} dK_R = \int_{\text{hela}} \frac{1}{2} \omega^2 dm r^2 = \frac{1}{2} \omega^2 \int_{\text{hela}} r^2 dm$$

$$\int r^2 dm = \text{tröghetsmomentet} = I$$

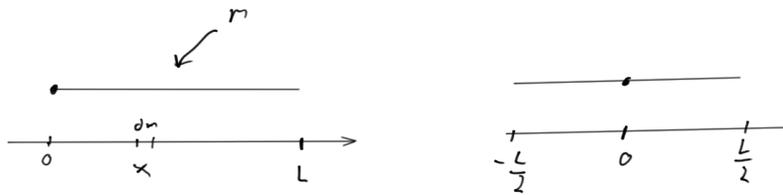
$$K_R = \frac{1}{2} I \omega^2 \text{ jämför med } K = \frac{1}{2} m v^2$$

## Exempel



## Smal pinne

$$\frac{dm}{M} = \frac{dx}{L} \rightarrow dm = \frac{M}{L} dx$$

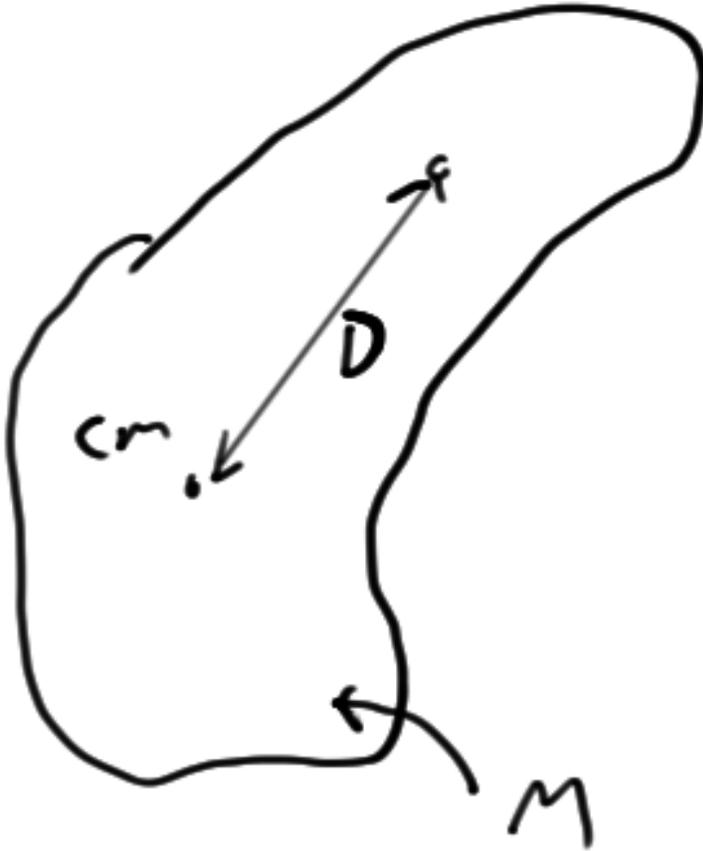


Fall 1:

$$I = \int_0^L x^2 dm = \int_0^L \frac{M}{L} x^2 dx = \frac{M}{L} \left[ \frac{1}{3} x^3 \right]_0^L = \frac{1}{3} ML^2$$

Fall 2:

$$I = \int_{-\frac{L}{2}}^{\frac{L}{2}} \frac{M}{L} x^2 dx = \frac{M}{L} \left[ \frac{1}{3} x^3 \right]_{-\frac{L}{2}}^{\frac{L}{2}} = \frac{ML^2}{12}$$

*Parallellförskjutningssatsen*

$$I_0 = I_{cm} + MD^2$$

Smal pinne:

$$I_{cm} = \frac{ML^2}{12}$$
$$I = \frac{ML^2}{12} + M \left( \frac{L}{2} \right)^2 = ML^2 \left( \frac{1+3}{12} \right) = \frac{ML^2}{3}$$

I för en homogen cylinder

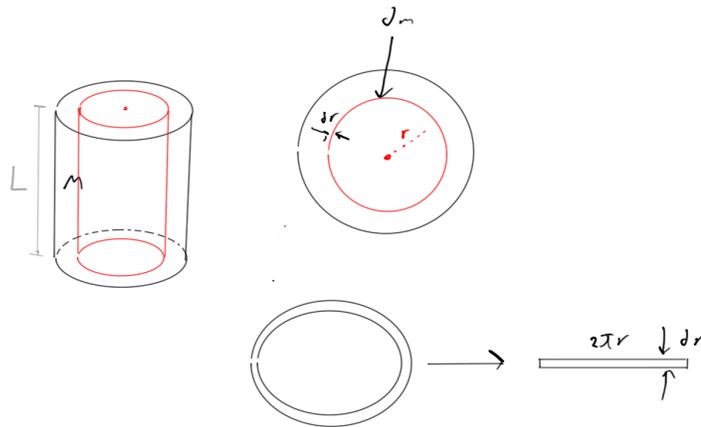
$$dI = r^2 dm$$

$$\rho = \frac{M}{\pi R^2 L} = \text{tätheten}$$

$$dV = (2\pi r) dr * L$$

$$dm = 2\pi r dr L \frac{M}{\pi R^2 L} = \frac{2M}{R^2} r dr$$

$$I = \int_0^R di = \frac{2M}{R^2} \int_0^R r * r^2 dr = \frac{2M}{R^2} \left[ \frac{R^4}{4} \right]_0^R = \frac{2M}{R^2} \frac{1}{4} R^4 = \frac{1}{2} MR^2$$



Kryssprodukt

$$\vec{A} \times \vec{B}$$

$$\hat{i} \times \hat{j} = \hat{k}$$

$$\hat{j} \times \hat{i} = -\hat{k}$$

$$|\vec{A} \times \vec{B}| = |\vec{A}| \cdot |\vec{B}| \cdot \sin \rho$$

Vridande moment

$$\vec{\tau} \equiv \vec{r} \times \vec{F}$$

## Rörelsemängdsmoment

$$\vec{L} \equiv \vec{r} \times \vec{p}$$

Samband mellan  $\vec{\tau}$  och  $\vec{L}$

$$\frac{d\vec{L}}{dt} = \frac{d}{dt}(\vec{r} \times m\vec{v}) = \frac{d\vec{r}}{dt} \times m\vec{v} + \vec{r} \times \frac{d}{dt}(m\vec{v}) = \vec{r} \times \vec{F} = \vec{\tau}$$

$$\vec{\tau} = \frac{d\vec{L}}{dt}$$